

The Pauli/Dirac Matrices

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This document shows the relationship of the Pauli/Dirac matrices to the relativistic unit circle

(I'm currently writing it, but this will give you a taste so far).

[The Relativistic Unit Circle](#) (under revision)

There are three Pauli Matrices for the case of two identical particles $\gamma=1$, $\beta=1$:

$$\sigma_1 = \begin{vmatrix} 0 & 1 \\ 1 & 0 \end{vmatrix} = \begin{vmatrix} 0 & \gamma \\ \gamma & 0 \end{vmatrix}, \gamma=1$$

$$\sigma_2 = \begin{vmatrix} 0 & i \\ -i & 0 \end{vmatrix} = \begin{vmatrix} 0 & i\beta \\ -i\beta & 0 \end{vmatrix}, \beta=1$$

$$\sigma_3 = \begin{vmatrix} 1^2 & 0 \\ 0 & i^2 \end{vmatrix} = \begin{vmatrix} \gamma & 0 \\ 0 & -\gamma^2 \end{vmatrix}, \gamma=1$$

$$\sigma_3 = \begin{vmatrix} 1 & 0 \\ 0 & -1 \end{vmatrix} = \begin{vmatrix} \gamma & 0 \\ 0 & -\gamma \end{vmatrix}, \gamma=1$$

$$\begin{vmatrix} \gamma & 0 \\ 0 & -\gamma \end{vmatrix} \vec{i} \otimes \vec{j} = -\gamma^2 \vec{k} = \gamma^2 (-\vec{k})$$

$$\begin{vmatrix} 0 & \gamma \\ -\gamma & 0 \end{vmatrix} \vec{j} \otimes \vec{i} = \gamma^2 \vec{k}$$

$$E_3 = (i \otimes j) + (j \otimes i) = (\gamma^2 + \gamma^2) [\vec{k} + (-\vec{k})] = 2(\gamma^2)(\vec{0}) = 2(1^2)(\vec{0})$$

Stern=Gerlach Experiment

Although the energy represented as null is equal and opposite, it is directed along an axis perpendicular to the direction of motion classically; globally, we call the direction of motion $\vec{v} \cdot \vec{I}$, and indicate the equal and opposite attractive forces by $(\vec{0} \cdot \vec{J})$

$$\sigma_1 + \sigma_2 = \begin{vmatrix} 0 & \gamma \\ \gamma & 0 \end{vmatrix} + \begin{vmatrix} 0 & i\beta \\ -i\beta & 0 \end{vmatrix} = \begin{vmatrix} 0 & \gamma + i\beta \\ \gamma - i\beta & 0 \end{vmatrix}$$

$$E_{\sigma_1 + \sigma_2} = \vec{i} \otimes \vec{j} = \begin{vmatrix} 0 & \gamma + i\beta \\ \gamma - i\beta & 0 \end{vmatrix} \vec{i} \otimes \vec{j} = -(\gamma + i\beta)(\gamma - i\beta) \vec{k} = (\gamma^2 + \beta^2)(-\vec{k})$$

$$\vec{j} \otimes \vec{i} = \begin{vmatrix} \gamma + i\beta & 0 \\ 0 & \gamma - i\beta \end{vmatrix} \vec{j} \otimes \vec{i} = (\gamma + i\beta)(\gamma - i\beta) \vec{k} = (\gamma^2 + \beta^2) \vec{k}$$

$$(\vec{i} \otimes \vec{j}) + (\vec{j} \otimes \vec{i}) = 2(\gamma^2 + \beta^2)(\vec{k} + (-\vec{k})) = 2(\gamma^2 + \beta^2)(\vec{0}) = (2(\gamma^2) + 2\beta^2)(\vec{0})$$

Here we see there is an extra term $2\beta^2$ that is different from the initial condition of σ_3

This term represents the energy carried to the system from the B field that was turned on to separate the "spins"

We can add this term to the σ_3 initial condition by assuming it is not interacting; thus the total system is represented by the interaction charges and a separate interacting B field:

$$\sigma_4 = \begin{vmatrix} \beta & 0 \\ 0 & -\beta \end{vmatrix}$$

$$\vec{i} \otimes \vec{j} = \begin{vmatrix} \beta & 0 \\ 0 & -\beta \end{vmatrix} \vec{i} \otimes \vec{j} = (\beta)(-\beta) \vec{k} = (\beta^2) \vec{k}$$

$$\vec{j} \otimes \vec{i} = \begin{vmatrix} 0 & \beta \\ -\beta & 0 \end{vmatrix} \vec{j} \otimes \vec{i} = (-\beta)(\beta) \vec{k} = (\beta^2) \vec{k}$$

$$(\vec{i} \otimes \vec{j}) + (\vec{j} \otimes \vec{i}) = 2\beta^2(\vec{0})$$

The total of the interactions from and σ_3 σ_4 as initial states is then

$E_3 + E_4 = 2\gamma^2(\bar{0}) + 2\beta^2(\bar{0}) = 2(\gamma^2 + \beta^2)$, and is equal to that of the final states:

$$E_{\sigma_1+\sigma_2} = E_{\sigma_3+\sigma_4} = 2(\gamma^2 + \beta^2)(\bar{0})$$

The Dirac γ^0 Matrix

The Dirac γ^0 matrix can then be represented from the relativistic unit circle as a ccw Lorentz rotation

corresponding to separate rotations of $\theta = \frac{\pi}{2}$ starting with $\theta = 0$:

$$\gamma^0 = \begin{vmatrix} 1^2 & 0 & 0 & 0 \\ 0 & 1^2 & 0 & 0 \\ 0 & 0 & -1^2 & 0 \\ 0 & 0 & 0 & -1^2 \end{vmatrix} = \begin{vmatrix} \gamma^2 & 0 & 0 & 0 \\ 0 & \beta^2 & 0 & 0 \\ 0 & 0 & (i\gamma)^2 & 0 \\ 0 & 0 & 0 & (i\beta)^2 \end{vmatrix} = \begin{vmatrix} \cos^2 \theta & 0 & 0 & 0 \\ 0 & \sin^2 \theta & 0 & 0 \\ 0 & 0 & -\cos^2 \theta & 0 \\ 0 & 0 & 0 & -\sin^2 \theta \end{vmatrix},$$

where the trace of the upper left quadrant together with the trace of the lower right quadrant yields the Pauli σ_3 matrix :

$$\sigma_3 = \begin{vmatrix} 1^2 & 0 \\ 0 & -1^2 \end{vmatrix} = \begin{vmatrix} Tr(A) & 0 \\ 0 & Tr(B) \end{vmatrix}, |A| = \begin{vmatrix} \cos^2 \theta & 0 \\ 0 & \sin^2 \theta \end{vmatrix}, |B| = \begin{vmatrix} -\cos^2 \theta & 0 \\ 0 & -\sin^2 \theta \end{vmatrix}$$

$$\gamma^0 = \begin{vmatrix} |A| & 0 \\ 0 & |B| \end{vmatrix}$$

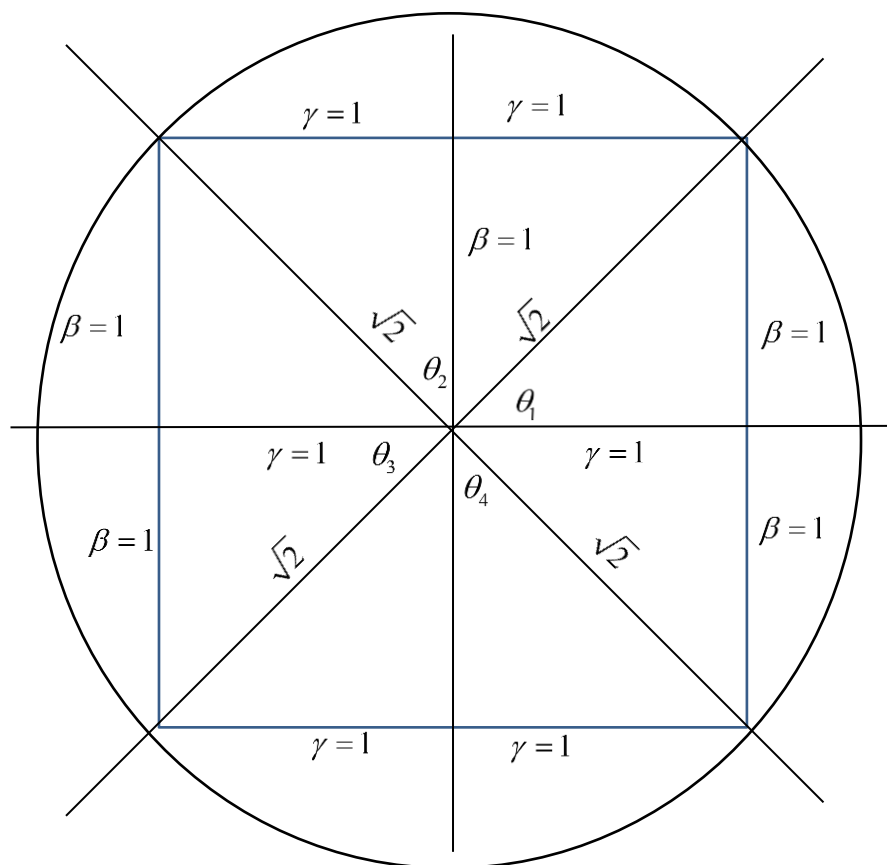
This matrix can be reconfigured to represent polarity reversal:

$$\gamma_0 = \begin{vmatrix} 1^2 & 0 & 0 & 0 \\ 0 & 1^2 & 0 & 0 \\ 0 & 0 & i^2 & 0 \\ 0 & 0 & 0 & i^2 \end{vmatrix} = \begin{vmatrix} \gamma & 0 & 0 & 0 \\ 0 & -\gamma & 0 & 0 \\ 0 & 0 & \beta & 0 \\ 0 & 0 & 0 & -\beta \end{vmatrix} = \begin{vmatrix} \cos \theta & 0 & 0 & 0 \\ 0 & -\cos \theta & 0 & 0 \\ 0 & 0 & \sin \theta & 0 \\ 0 & 0 & 0 & -\sin \theta \end{vmatrix}$$

Spin and Angular Momentum quantization is represented by Lorentz rotation to the diagonals, where

$\gamma = \beta$ introduces a factor of $\frac{1}{\sqrt{2}}$, so that an initial Lorentz rotation of $\frac{\pi}{4}$ corresponds to:

$$\gamma_0 = \begin{vmatrix} \frac{1}{\sqrt{2}}\gamma & 0 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}}\beta & 0 & 0 \\ 0 & 0 & -\frac{1}{\sqrt{2}}\gamma & 0 \\ 0 & 0 & 0 & -\frac{1}{\sqrt{2}}\beta \end{vmatrix} = \begin{vmatrix} \frac{1}{\sqrt{2}}\cos\theta & 0 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}}\sin\theta & 0 & 0 \\ 0 & 0 & -\frac{1}{\sqrt{2}}\cos\theta & 0 \\ 0 & 0 & 0 & -\frac{1}{\sqrt{2}}\sin\theta \end{vmatrix}$$



CCW Rotation

Positive Definite $(\gamma^0)^2$

$$(\gamma^0)^2 = I^4 = \begin{vmatrix} 1^4 & 0 & 0 & 0 \\ 0 & 1^4 & 0 & 0 \\ 0 & 0 & 1^4 & 0 \\ 0 & 0 & 0 & 1^4 \end{vmatrix} = \begin{vmatrix} \gamma^4 & 0 & 0 & 0 \\ 0 & \beta^4 & 0 & 0 \\ 0 & 0 & (i\gamma)^4 & 0 \\ 0 & 0 & 0 & (i\beta)^4 \end{vmatrix} = \begin{vmatrix} \cos^4 \theta & 0 & 0 & 0 \\ 0 & \sin^4 \theta & 0 & 0 \\ 0 & 0 & \cos^4 \theta & 0 \\ 0 & 0 & 0 & \sin^4 \theta \end{vmatrix}$$

Note that the trace of both matrices in the upper left and lower right quadrants is now:

$$\text{Tr}(A) = \text{Tr}(B) = \cos^4 \theta + \sin^4 \theta$$

Compare this result with the equation:

$$(\cos^2 \theta + \sin^2 \theta)^2 = \cos^4 \theta + 2\sin^2 \theta \cos^2 \theta + \sin^4 \theta = \gamma^4 + 2\gamma^2 \beta^2 + \beta^4$$

Note that there is now an "interaction" term $2\gamma^2 \beta^2$ that now must be accounted for in the expansion for the positive definite representation of γ^0 . This is done by the remaining Dirac gamma matrices, γ^1 , γ^2 and γ^3 .

$$\gamma^0 = (\sqrt{\gamma^0})^2 = \begin{vmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & 0 & 0 & i \end{vmatrix} \begin{vmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & 0 & 0 & i \end{vmatrix} = \begin{vmatrix} 1^2 & 0 & 0 & 0 \\ 0 & 1^2 & 0 & 0 \\ 0 & 0 & i^2 & 0 \\ 0 & 0 & 0 & i^2 \end{vmatrix} = \begin{vmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{vmatrix}$$

$$(\gamma^0)^2 = \begin{vmatrix} 1^2 & 0 & 0 & 0 \\ 0 & 1^2 & 0 & 0 \\ 0 & 0 & i^2 & 0 \\ 0 & 0 & 0 & i^2 \end{vmatrix} = \begin{vmatrix} 1^4 & 0 & 0 & 0 \\ 0 & 1^4 & 0 & 0 \\ 0 & 0 & 1^4 & 0 \\ 0 & 0 & 0 & 1^4 \end{vmatrix}$$

$$\gamma^0 = (\sqrt{\gamma^0})^2 = \begin{vmatrix} \gamma & 0 & 0 & 0 \\ 0 & \gamma & 0 & 0 \\ 0 & 0 & i\beta & 0 \\ 0 & 0 & 0 & i\beta \end{vmatrix} \begin{vmatrix} \gamma & 0 & 0 & 0 \\ 0 & \gamma & 0 & 0 \\ 0 & 0 & i\beta & 0 \\ 0 & 0 & 0 & i\beta \end{vmatrix} = \begin{vmatrix} \gamma^2 & 0 & 0 & 0 \\ 0 & \gamma^2 & 0 & 0 \\ 0 & 0 & -\beta^2 & 0 \\ 0 & 0 & 0 & -\beta^2 \end{vmatrix}$$

$$\gamma^1 = \begin{vmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & 0 & 0 & i \end{vmatrix} \begin{vmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & i & 0 & 0 \\ i & 0 & 0 & 0 \end{vmatrix} = \begin{vmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{vmatrix}$$

$$= \left[(\sqrt{\gamma_0})(\sqrt{\gamma_0})^T \right]$$

$$\gamma^2 = \begin{vmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ -i & 0 & 0 & 0 \end{vmatrix}$$

$$\gamma^3 = \begin{vmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{vmatrix} = \begin{vmatrix} 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{vmatrix} \begin{vmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{vmatrix}$$

$$(\gamma^0)^2 = \begin{vmatrix} 1^2 & 0 & 0 & 0 \\ 0 & 1^2 & 0 & 0 \\ 0 & 0 & i^2 & 0 \\ 0 & 0 & 0 & i^2 \end{vmatrix} = \begin{vmatrix} 1^4 & 0 & 0 & 0 \\ 0 & 1^4 & 0 & 0 \\ 0 & 0 & 1^4 & 0 \\ 0 & 0 & 0 & 1^4 \end{vmatrix}$$

$$(\gamma^1)^2 = \begin{vmatrix} 0 & 0 & 0 & 1^2 \\ 0 & 0 & 1^2 & 0 \\ 0 & i^2 & 0 & 0 \\ i^2 & 0 & 0 & 0 \end{vmatrix} = \begin{vmatrix} -(1^4) & 0 & 0 & 0 \\ 0 & -(1^4) & 0 & 0 \\ 0 & 0 & -(1^4) & 0 \\ 0 & 0 & 0 & -(1^4) \end{vmatrix}$$

$$(\gamma^2)^2 = \begin{vmatrix} 0 & 0 & 0 & -i^2 \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ -i & 0 & 0 & 0 \end{vmatrix} = \begin{vmatrix} -(1)^2 & 0 & 0 & 0 \\ 0 & -(1)^2 & 0 & 0 \\ 0 & 0 & -(1)^2 & 0 \\ 0 & 0 & 0 & -(1)^2 \end{vmatrix} = \begin{vmatrix} i^2 & 0 & 0 & 0 \\ 0 & i^2 & 0 & 0 \\ 0 & 0 & i^2 & 0 \\ 0 & 0 & 0 & i^2 \end{vmatrix}$$

$$(\gamma^3)^2 = \begin{vmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{vmatrix}^2 = \begin{vmatrix} 0 & 0 & \gamma & 0 \\ 0 & 0 & 0 & (i\gamma) \\ (i\gamma) & 0 & 0 & 0 \\ 0 & \gamma & 0 & 0 \end{vmatrix}^2 = \begin{vmatrix} -1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{vmatrix} = \begin{vmatrix} i^2 & 0 & 0 & 0 \\ 0 & i^2 & 0 & 0 \\ 0 & 0 & i^2 & 0 \\ 0 & 0 & 0 & i^2 \end{vmatrix}$$

$$\gamma^5 = \begin{vmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{vmatrix}$$

$$\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$$

The interaction product

The Interaction Product is the area ("energy") of the interaction triangle

$A_{\gamma\beta} = \frac{1}{2}\gamma\beta = \frac{1}{2}1^2$ in each quadrant at the null vector at the vector connection (or at the circumference only if equal and opposite as the tangent, since independent on the plane (γ, β)) For $\gamma = \beta = 1$, there is a phase difference of $\Delta\theta = \frac{\pi}{2}$ in each quadrant, and of $\Delta\theta = \pi$ for each axis considered separately.

For each axis separately, with equal and opposite interaction at the origin:

$(+\gamma, -\gamma)$, $(+\beta, -\beta)$ phase difference π

$$\begin{vmatrix} \beta & 0 \\ 0 & -\beta \end{vmatrix} \rightarrow (1^2)\beta^2(\vec{0})$$

$$\begin{vmatrix} \gamma & 0 \\ 0 & -\gamma \end{vmatrix} \rightarrow (1^2)\gamma^2(\vec{0})$$

For $\gamma = \beta = 1$, the sum is equal to $2\gamma\beta(\vec{0}) = 2(1^2)(\vec{0})$

Energy as sum of quadrants

For the energy in each quadrant, Lorentz rotating ccw in the relativistic unit circle with phase difference $\frac{\pi}{2}$:

$(+\gamma, +\beta)$ 1st Quadrant

$$\begin{vmatrix} 1 & 0 \\ 0 & 1 \end{vmatrix} \rightarrow (\vec{i} \otimes \vec{j}) = 1^2(\vec{k})$$

$$\begin{vmatrix} 0 & 1 \\ 1 & 0 \end{vmatrix} \rightarrow (\vec{j} \otimes \vec{i}) = 1^2(\vec{-k})$$

$$(\vec{i} \otimes \vec{j}) + (\vec{j} \otimes \vec{i}) = 1^2(\vec{0})$$

Interaction Product:

$$\begin{vmatrix} \gamma & 0 \\ 0 & \beta \end{vmatrix} \rightarrow \left(\frac{1}{2}\right)\beta\gamma(\vec{0})$$

$(+\beta, -\gamma)$ 2nd quadrant

$$\begin{vmatrix} 1 & 0 \\ 0 & -1 \end{vmatrix} \rightarrow (\vec{i} \otimes \vec{j}) = 1^2 (\vec{-k})$$

$$\begin{vmatrix} 0 & 1 \\ -1 & 0 \end{vmatrix} \rightarrow (\vec{j} \otimes \vec{i}) = 1^2 (\vec{k})$$

$$(\vec{i} \otimes \vec{j}) + (\vec{j} \otimes \vec{i}) = 1^2 (\vec{0})$$

Interaction Product:

$$\begin{vmatrix} \beta & 0 \\ 0 & -\gamma \end{vmatrix} \rightarrow \left(\frac{1}{2}\right) \beta \gamma (\vec{0})$$

$(-\gamma, -\beta)$ 3rd quadrant

$$\begin{vmatrix} -1 & 0 \\ 0 & -1 \end{vmatrix} \rightarrow (\vec{i} \otimes \vec{j}) = 1^2 (\vec{k})$$

$$\begin{vmatrix} 0 & -1 \\ -1 & 0 \end{vmatrix} \rightarrow (\vec{j} \otimes \vec{i}) = 1^2 (\vec{-k})$$

$$(\vec{i} \otimes \vec{j}) + (\vec{j} \otimes \vec{i}) = 1^2 (\vec{0})$$

Interaction Product

$$\left(\frac{1}{2}\right) \begin{vmatrix} -\beta & 0 \\ 0 & -\gamma \end{vmatrix} \rightarrow \left(\frac{1}{2}\right) \beta \gamma (\vec{0})$$

$(-\beta, +\gamma)$ 4th quadrant

$$\begin{vmatrix} -1 & 0 \\ 0 & 1 \end{vmatrix} \rightarrow (\vec{i} \otimes \vec{j}) = 1^2 (\vec{-k})$$

$$\begin{vmatrix} 0 & -1 \\ 1 & 0 \end{vmatrix} \rightarrow (\vec{j} \otimes \vec{i}) = 1^2 (\vec{k})$$

$$(\vec{i} \otimes \vec{j}) + (\vec{j} \otimes \vec{i}) = 1^2 (\vec{0})$$

Interaction Product:

$$\left(\frac{1}{2}\right) \begin{vmatrix} -\beta & 0 \\ 0 & \gamma \end{vmatrix} \rightarrow \left(\frac{1}{2}\right) \beta \gamma (\vec{0})$$

$$E(\vec{0}) = 4 \left(\frac{1}{2} \gamma \beta\right) = 2\gamma\beta$$

The interaction energy can also be calculated for $\gamma = \beta$ in each quadrant as the matrix (e.g. in the first quadrant) as

$$\begin{vmatrix} \frac{\gamma}{\sqrt{2}} & 0 \\ 0 & \frac{\beta}{\sqrt{2}} \end{vmatrix} \rightarrow \frac{1}{2}(\gamma\beta), \text{ consistent with the "argument by geometry" above.}$$

This interaction relation is actually spin polarization energy of photon on photon interaction (with β representing a change (or perturbation) to the initial "rest energy of the photon γ ". However, if β is characterized as the "spin" of a photon-equivalent electron (i.e., a massive heavier photon), then the "spin" is independent of the electron mass as defined by the relativistic unit circle; if there is no "spin" interaction with an external field, then the initial state of the electron is the same as the final state $\alpha = \alpha'$, $\gamma^2 = 1^2$, $\beta^2 = 0$ and its "rest energy" or "rest mass" is invariant.

The spin can be quantized; for the case $\gamma = \beta = 1 = h$ describing two photon - equivalent electrons (or photons) interaction, the interaction energy is then defined as $E_s = s^2 = \frac{1}{2}h^2$ for each quadrant of the relativistic circle, so that

$$s = \frac{1}{\sqrt{2}}h = \frac{h}{\sqrt{2}}. \text{ For } n \text{ such interactions, then } s_n = \frac{nh}{\sqrt{2}} = ns$$

Consider the case where $\gamma\beta > 1$ (i.e., the initial state is invariant). Then

$$(\gamma')^2 = 1^2 + (\gamma\beta)^2 \text{ Since } \gamma = 1 \text{ in the case of the relativistic unit circle, this becomes}$$

$$(\gamma')^2 = \gamma^2 + (\gamma\beta)^2, \text{ where } \gamma \text{ can be thought of as the photo-equivalent electron.}$$

Then for n such electrons, the equation becomes

$$(\gamma')^2 = (n\gamma)^2 + (n\gamma)^2(\beta)^2 = n^2 + n(\beta)^3 \text{ However, since } \beta^3 = (1)^3, \beta > 0 \text{ the equation can}$$

be written as $(\gamma')^2 = n^2 + n(1)^3 = n(n + (1)^3)$, (where $(1)^3$ usually appears as "1", so

that $(\gamma')^2 = n^2 + n$), bearing in mind that β is always positive if the interaction is

between bases in the first quadrant. Note that if $\beta = 0$, $(\gamma')^2 = n^2$ (i.e., there is no

interaction). In particular, $\lim_{n \rightarrow \infty} (\gamma')^2 = \lim_{n \rightarrow \infty} n^2 + n = \lim_{n \rightarrow \infty} n(n+1) = n^2$; i.e., n^2 is analogous

to a "continuum" in the state model of atomic physics (note that β^3 suggests a

three dimensional "volume", but the precise discussion is outside the scope of this document)..

From the Binomial Expansion, $\psi^n = (\gamma + \beta)^n = \gamma^n + \beta^n + \text{rem}(\gamma, \beta, n)$ For the case $n = 2$, $\text{rem}(\gamma, \beta, 2) = 2\gamma\beta$ which is always positive in the first quadrant (and the third), so that it only vanishes if $\beta = 0, \gamma = 1$ (there is only one non-interacting photon or photo-equivalent electron).